Chapter 1. Heat Equation

Applied Partial Differential Equations with Fourier Series and Boundary Value Problems 5th Edition Richard Haberman Solutions Manual

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- 1.2.9 (d) Circular cross section means that P = 2πr, A = πr², and thus P/A = 2/r, where r is the radius. Also γ = 0.
- 1.2.9 (e) u(x, t) = u(t) implies that

$$c\rho \frac{d\omega}{dr} = -\frac{2h}{r}\omega$$
.

The solution of this first-order linear differential equation with constant coefficients, which satisfies the initial condition $u(0) = u_{0}$, is

$$m(r) = m_0 \exp -\frac{2h}{c\rho r}r^*.$$

Section 1.3

1.3.2 $\partial w/\partial x$ is continuous if $K_{\alpha}(x_{\alpha}-) = K_{\alpha}(x_{\alpha}+)$, that is, if the conductivity is continuous.

Section 1.4

- 14.1 (a) Equilibrium satisfies (1.4.14), d²u/dx² = 0, whose general solution is (1.4.17), u = c₁ + c₂x. The boundary condition u(0) = 0 implies c₁ = 0 and u(L) = T implies c₂ = T/L so that u = Tx/L.
- 1.4.1 (d) Equilibrium satisfies (1.4.14), d²u/dx² = 0, whose general solution (1.4.17), u = c₁ + c₂x. From the boundary conditions, u(0) = T yields T = c₁ and du/dx(L) = α yields α = c₂. Thus u = T + αx.
- 1.4.1 (f) In equilibrium, (1.2.9) becomes d²u/dx² = -Q/K₀ = -x², whose general solution (by integrating twice) is u = -x⁴/12 + c₁ + c₂x. The boundary condition u(0) = T yields c₁ = T, while du/dx(L) = 0 yields c₂ = L³/3. Thus u = -x⁴/12 + L³x/3 + T.
- 1.4.1 (h) Equilibrium satisfies $d^2u/dx^2 = 0$. One integration yields $du/dx = c_2$, the second integration yields the general solution $u = c_1 + c_2 x$.

$$x = 0$$
: $c_2 - (c_1 - T) = 0$
 $x = L$: $c_2 = a$ and thus $c_1 = T + a$.

Therefore, $u = (T + \alpha) + \alpha x = T + \alpha (x + 1)$.

1.4.7 (a) For equilibrium:

$$\frac{d^2u}{dx^2} = -1$$
 implies $u = -\frac{x^2}{2} + c_1x + c_2$ and $\frac{du}{dx} = -x + c_1$.

From the boundary conditions $\frac{\partial G}{\partial x}(0) = 1$ and $\frac{\partial G}{\partial x}(L) = \beta$, $c_1 = 1$ and $-L + c_1 = \beta$ which is consistent only if $\beta + L = 1$. If $\beta = 1 - L$, there is an equilibrium solution ($\mu = -\frac{2}{2} + x + c_2$). If $\beta = 1 - L$, there isn't an equilibrium solution. The difficulty is caused by the heat flow being specified at both ends and a source specified inside. An equilibrium will exist only if these three are in balance. This

Partial Differential Equations And Boundary Value <u>Problems</u>

Stephen A. Wirkus,Randall J. Swift,Ryan Szypowski

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Partial Differential Equations and Boundary-Value Problems with Applications Mark A. Pinsky,2011 Building on the basic techniques of separation of variables and Fourier series the book presents the solution of boundary value problems for basic partial differential equations the heat equation wave equation and Laplace equation considered in various standard coordinate systems rectangular cylindrical and spherical Each of the equations is derived in the three dimensional context the solutions are organized according to the geometry of the coordinate system which makes the mathematics especially transparent Bessel and Legendre functions are studied and used whenever appropriate throughout the text The notions of steady state solution of closely related stationary solutions are developed for the heat equation applications to the study of heat flow in the earth are presented The problem of the vibrating string is studied in detail both in the Fourier transform setting and from the viewpoint of the explicit representation d Alembert formula Additional chapters include the numerical analysis of solutions and the method of Green s functions for solutions of partial differential equations The exposition also includes asymptotic methods Laplace transform and stationary phase With more than 200 working examples and 700 exercises more than 450 with answers the book is suitable for an undergraduate course in partial differential equations

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